Incoherent Dicke model - new type of phase transition

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# **Motivation 1**

• Deep strong coupling  $\rightarrow$  Rabi ground state changes qualitatively.

 $\bullet$  Deep strong coupling  $\rightarrow$  Dicke 'superradiant' phase transition.

✦ Hard to reach, classical Dicke

$$\hat{H} = \kappa \hat{S}_z + \omega \hat{S}_x$$

Not critical, not even very exciting!

'Incoherent Dicke model'

$$\partial \hat{\rho} = i \left[ \hat{\rho}, \omega \hat{S}_x \right] + \kappa \left( 2\hat{S}^- \hat{\rho}\hat{S}^+ - \hat{S}^+ \hat{S}^- \hat{\rho} - \hat{\rho}\hat{S}^+ \hat{S}^- \right)$$

It's getting interesting....

# **Motivation 2**

Equilibrium phase transitions (continuous):

- **Universal** few critical exponents determine the physics (microscopic details irrelevant).
- **Spontaneous symmetry breaking** in one phase the state does not possess the same symmetry as the full Hamiltonian.
- **Excitations** energy gap closes at the critical point. Excitations in 'symmetry broken phase' either gapped (*Higgs*) or continuous (*Goldstone*).
- *Mermin-Wagner theorem* the type of symmetry + the dimensionality determine which type of transition that is allowed.

How do the above results translate to phase transitions in open quantum systems?

# 1. (Equilibrium) Quantum phase transitions

- Symmetry breaking.
- Scale invariance and universality.
- 2. Open quantum phase transitions
- 3. A new type of phase transition
- 4. Prospects

# 1. (Equilibrium) Quantum phase transitions

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# Phase transitions Phase transition

← Action  $S[\bar{\psi}, \psi]$ , giving partition function

$$\mathcal{Z} = \int D(\bar{\psi}, \psi) \, e^{-S[\bar{\psi}, \psi]}$$

Mean-field solution: minimizing the action for order parameter  $\psi$ .

• Symmetries  $S[g\bar{\psi},g\psi] = S[\bar{\psi},\psi]$ .



# Phase transitions 2nd order phase transition

- ✦ Continuous phase transition.
- Exists a symmetry  $[\hat{U}, \hat{H}] = 0$ .
- ✦ Quantum mechanics: Û and Ĥ common eigenbasis energy eigenstates well defined symmetry.
- ← Thermodynamic limit: *spontaneous symmetry breaking* ground state  $|\psi_0(\lambda)\rangle$  does not have a defined symmetry!
- $\blacklozenge$  "Potential barrier" becomes infinite  $\rightarrow$  degeneracy.
- ♦ If Û continuous → Goldstone (gapless) excitations.
  Û discontinuous → Higgs (gapped) excitations.

N. Goldenfeld, Lectures on Phase Transitions...

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## **Phase transitions** *Phase transition (2nd) - scale invariance*

At the critical point  $\lambda_c$ 

The following Ising model configurations range from 2048x2048 sites to 131072x131072 sites.

$$H_{class} = \sum_{\langle ij \rangle} s_i s_j - h \sum_i s_i$$

Can you tell which is which?

System the same independent of "zooming" - length scale diverges!!

N. Goldenfeld, Lectures on Phase Transitions...

# **Phase transitions** *Phase transition (2nd) - universality*

 Systematic method (*renormalization group*) - eliminates short length scales ('high energies').

✦ Effective low energy model - microscopic details irrelevant.

 Models belong to different 'classes' - universality (depend on macroscopic properties like symmetries).

Critical regime:

 $\xi \propto |\lambda - \lambda_c|^{-\nu}$ , Characteristic length  $\tau \propto |\lambda - \lambda_c|^{-\delta}$ , Characteristic time  $\Delta E \propto |\lambda - \lambda_c|^{z\nu}$ , Gap closing

*Critical exponents*  $\nu$ ,  $\delta$ , z,..., different *universality classes*.

N. Goldenfeld, Lectures on Phase Transitions...

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### Phase transitions in open quantum systems

- Quantum PT's, non-analyticity in ground state  $|\psi_0(\lambda)\rangle$  at critical coupling  $\lambda_c$ .
- Transition due to quantum fluctuations.
- ✦ Open quantum system

$$\partial_t \hat{\rho}(t) = i \left[ \hat{\rho}(t), \hat{H}_S \right] + \hat{\mathcal{L}} \left[ \hat{\rho}(t) \right],$$
$$\hat{\mathcal{L}} \left[ \hat{\rho}(t) \right] = \sum_i g_i \left( 2\hat{A}_i \hat{\rho}(t) \hat{A}_i^{\dagger} - \hat{A}_i^{\dagger} \hat{A}_i \hat{\rho}(t) - \hat{\rho}(t) \hat{A}_i^{\dagger} \hat{A}_i \right)$$

✦ No ground state, no energy spectrum!!

C.-E. Bardyn *et al.*, NJP **15** (2013)

### Phase transitions in open quantum systems

One (obvious) physically relevant state is the steady state

$$\partial_t \hat{\rho}_{ss}(t) = 0 \quad \rightarrow \quad i \left[ \hat{\rho}_{ss}, \hat{H}_S \right] + \hat{\mathcal{L}}[\hat{\rho}_{ss}] = 0$$

May be non-equilibrium, but time-independent.

Closed system, all energy eigenstates also steady states.

#### ✦ If:

- 1.  $[\hat{A}_i, \hat{H}_S] = 0, \ \forall \hat{A}_i$ , energy eigenstates also steady states.
- 2.  $\hat{H}_S$  critical and  $\hat{\mathcal{L}}[|\psi_0(\lambda < \lambda_c)\rangle\langle\psi_0(\lambda < \lambda_c)|] = 0$ , the environment is expected to support the symmetric phase.
- If neither of the two  $\rightarrow$  new physics???
- $\blacklozenge$  Phase transition if  $\hat{\rho}_{ss}$  non-analytic for some  $\lambda_c$  .

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  - Composite systems, density operators.
  - Lindblad master equation.
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# 4. A new type of phase transition

## 5. Prospects

## New type of phase transition Model

✦ Simplest (non-trivial) scenario

- $\partial_t \hat{\rho} = i[\hat{\rho}, \hat{H}_s]$  trivial
- $\partial_t \hat{\rho} = \kappa (2\hat{A}\hat{\rho}\hat{A}^{\dagger} \hat{A}^{\dagger}\hat{A}\hat{\rho} \hat{\rho}\hat{A}^{\dagger}\hat{A})$  trivial
- $\partial_t \hat{\rho} = i[\hat{\rho}, \hat{H}_s] + \kappa (2\hat{A}\hat{\rho}\hat{A}^{\dagger} \hat{A}^{\dagger}\hat{A}\hat{\rho} \hat{\rho}\hat{A}^{\dagger}\hat{A})$  critical

One such model (which is also solvable analytically!)

 $[\hat{S}_i, \hat{S}_j] = i \varepsilon_{ijk} \hat{S}_k$  Large spin-S (preserved)

 $\hat{H}_S = \omega \hat{S}_x$  Coherent drive

 $\hat{A} = \hat{S}_{-}$  Spontaneous decay to state  $|S, -S\rangle$ 

D. F. Walls et al., Prog. Theo. Phys. 64 (1978)

## New type of phase transition Model

✦ Equation to solve

$$\partial \hat{\rho} = i[\hat{\rho}, \omega \hat{S}_x] + \kappa (2\hat{S}_-\hat{\rho}\hat{S}_+ - \hat{S}_+\hat{S}_-\hat{\rho} - \hat{\rho}\hat{S}_+\hat{S}_-)$$

✦ Limiting cases

$$\frac{\kappa}{\omega} = 0 \quad \rightarrow \quad |S, -S\rangle_x$$
 Hamiltonian dominates

$$\frac{\omega}{\kappa}=0 ~~ 
ightarrow~ |S,-S
angle_z$$
 Lindbladian dominates

Phase transition somewhere between?

### New type of phase transition Mean-field solution

• Mean-field (normal ordering):  $S_{\alpha} = \langle \hat{S}_{\alpha} \rangle, \quad \langle \hat{S}_{\alpha} \hat{S}_{\beta} \rangle = \langle \hat{S}_{\alpha} \rangle \langle \hat{S}_{\beta} \rangle.$ 

$$\dot{S}_x = 2\frac{\kappa}{S}S_x S_z,$$

$$\dot{S}_y = S_z \left( 2\frac{\kappa}{S} S_y - \omega \right),$$

$$\dot{S}_z = \omega S_y - 2\frac{\kappa}{S} \left(S_x^2 + S_y^2\right)$$

♦ Conserved spin:  $S^2 = S_x^2 + S_y^2 + S_z^2$ .

Steady state solutions

$$(S_x, S_y, S_z)/S = \left(\pm\sqrt{1 - \frac{4\kappa^2}{\omega^2}}, \frac{2\kappa}{\omega}, 0\right), \qquad (S_x, S_y, S_z)/S = \left(0\frac{\omega}{2\kappa}, \pm\sqrt{1 - \frac{\omega^2}{4\kappa^2}}\right)$$

## New type of phase transition Mean-field solution



# New type of phase transition Universality

✦ Universality - critical exponents.

- Mean-field.  $\lambda = \omega/\kappa \rightarrow$  "Magnetization"  $S_z \propto |\lambda_c \lambda|^{1/2} \rightarrow \nu = 1/2$ .
- Quantum.  $\nu = 1/2$ .
- ✦ Seems to be universal!
- No length scale! Critical slowing down??

# New type of phase transition Universality

✦ Quantum treatment. Analytical solution

$$\hat{\rho}_{ss} = \frac{1}{D} \sum_{n,m=0}^{2S} (g^*)^{-m} g^{-n} \hat{S}^m_- \hat{S}^n_+, \qquad g = i\omega S/\kappa$$

Expectations  $\langle \hat{\mathcal{O}} \rangle = \text{Tr}[\hat{\mathcal{O}}\hat{\rho}_{ss}]$ . Truncate for some S.

- ✦ Critical exponent the same as for mean-field.
- Fluctuations in  $\hat{S}_x$ :  $\langle \hat{S}_x^2 \rangle \propto |\lambda_c \lambda|^{1/2}$ .

# New type of phase transition Symmetries

#### ✦ Symmetries.

- Closed case.  $[\hat{A}, \hat{H}] = 0 \leftrightarrow \hat{H} = \hat{U}\hat{H}\hat{U}^{-1} \rightarrow \hat{A}$  conserved.
- Open case.  $\partial \hat{\rho} = \mathcal{L}_{\hat{L}}[\hat{\rho}] = \mathcal{L}_{\hat{U}\hat{L}\hat{U}^{-1}}[\hat{\rho}] \longrightarrow \text{generally no conserved quantity.}$

◆ Guess 
$$\hat{S}_x \rightarrow -\hat{S}_x, \ \hat{S}_y \rightarrow \hat{S}_y, \ \hat{S}_z \rightarrow -\hat{S}_z$$

$$\partial \hat{\rho} = i[\hat{\rho}, \omega \hat{S}_x] + \kappa \left( 2\hat{S}_- \hat{\rho}\hat{S}_+ - \hat{S}_+ \hat{S}_- \hat{\rho} - \hat{\rho}\hat{S}_+ \hat{S}_- \right) \\
\rightarrow \\
\partial \hat{\rho} = -i[\hat{\rho}, \omega \hat{S}_x] + \kappa \left( 2\hat{S}_+ \hat{\rho}\hat{S}_- - \hat{S}_- \hat{S}_+ \hat{\rho} - \hat{\rho}\hat{S}_- \hat{S}_+ \right)$$

Dual symmetry! Flip spectrum + "step up" instead of "step down".

# New type of phase transition Symmetries

### ✦ Mean-field.

- First bifurcation (Sz). No pitchfork, only one or two solutions. Stability: one stable, one unstable.
- Second bifurcation (S<sub>x</sub>). Stability: the two solutions with purely imaginary eigenvalues → Hopf. (Phase space a sphere so peculiar trajectories).



# New type of phase transition

Second order phase transition with no symmetry breaking!!

- ✦ Is it a quantum phase transition?
  - Single degree of freedom. Quantum fluctuations vanishingly small in thermodynamic limit.
  - Fluctuations due to coupling to reservoir.
  - Quantum critical region signatures of "quantum" survives finite temperatures in the vicinity of a critical point.



S. Sachdev, Quantum Phase Transitions

#### 

- Coherence: *Purity*  $P \equiv \text{Tr}[\hat{\rho}^2] > 0$  (thermodynamic limit).
- Entanglement: Concurrence  $C(\hat{\rho})$  between two qubits. Negativity  $N(\hat{\rho})$  between two equal halves.



Thermodynamic limit:

Reservoir dominated - pure state.

Hamiltonian dominated - maximally mixed.

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- Coherence: *Purity*  $P \equiv \text{Tr}[\hat{\rho}^2] > 0$  (thermodynamic limit).
- Entanglement: *Concurrence*  $C(\hat{\rho})$  between two qubits. *Negativity*  $N(\hat{\rho})$  between two equal halves.



Thermodynamic limit:

No qubit-qubit entanglement outside the critical point.

#### 

- Coherence: *Purity*  $P \equiv \text{Tr}[\hat{\rho}^2] > 0$  (thermodynamic limit).
- Entanglement: Concurrence  $C(\hat{\rho})$  between two qubits. Negativity  $N(\hat{\rho})$  between two equal halves.



Thermodynamic limit:

No subsystem-subsystem entanglement outside the critical point.

Scales  $N(\hat{\rho}) \sim S$  ('volume law').

'Critical exponents' infinite.

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### Phase transitions in open quantum systems

✦ Found a model possessing an open PT, seemingly second order but lacks any symmetry breaking → New type. Partly unexplored field.

#### **Open questions**

- Are these PT's universal? Scale invariance? New classes?
- Models driven more strongly by quantum fluctuations, length scale (lattice models).
- Gap closening of the *Liouvillian*, universal?
- Is there a *Mermin-Wagner theorem* for open QPT's?
- *Kibble-Zurek mechanism*?

# Thanks!