

THE FUTURE OF BOSONIZATION

(1)

1. HISTORY

2. DIMENSIONS

ZERO : λ -RAY

ONE : HUTTINGER, HUBBARD

TWO : SQUARE FERMI SURFACE

THREE : QED (QCD?)

3. WHY?

λ -RAY THRESHOLD PROBLEM

SCHOTTE & SCHOTTE

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TOMONAGA, HUTTINGER, LIEB & MATTIS
 (INTERACTING ELECTRONS, SPINLESS, IN
 ONE SPACE DIMENSION)

$$H = \sum_k (k - k_F) a_{1k}^\dagger a_{1k} + (-k - k_F) a_{2k}^\dagger a_{2k} \\
 + \frac{g}{L} \sum_q P_1(q) P_2(-q)$$

a_{1k}, a_{2k} — FERMIONS

$$[a_{1k}, a_{1k'}^\dagger] = \delta_{kk'}$$

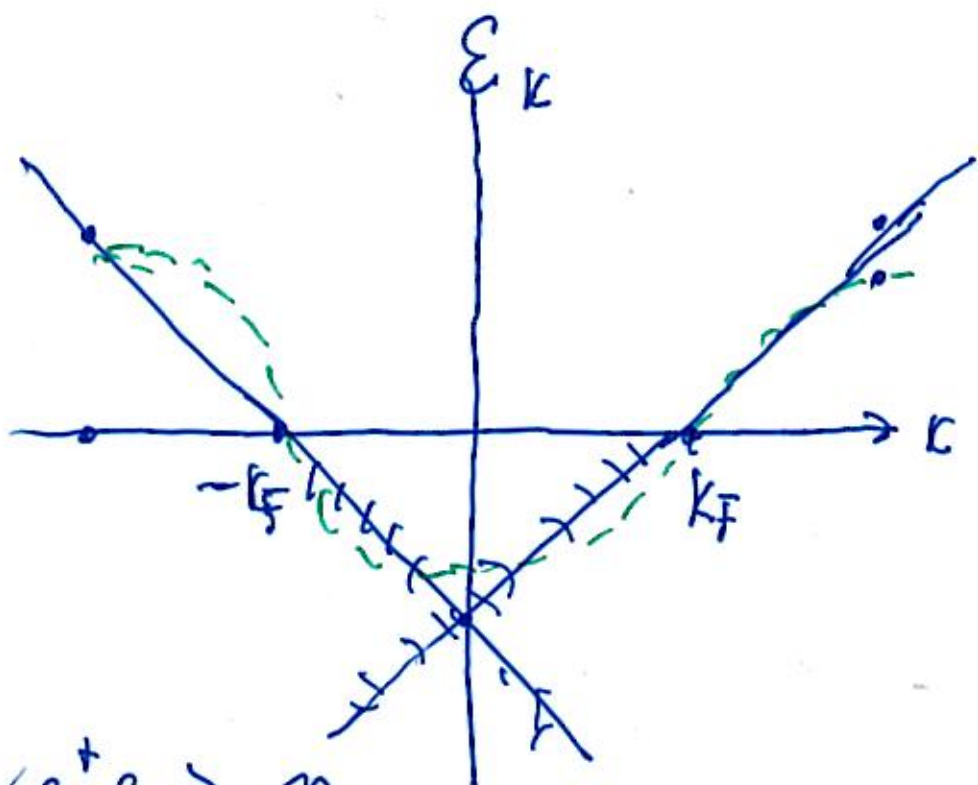
$$P_1(q) = \sum_k a_{1k+q}^\dagger a_{1k} \quad \text{etc.}$$

$$k = \frac{2\pi D}{L} \quad (n \text{ integer})$$

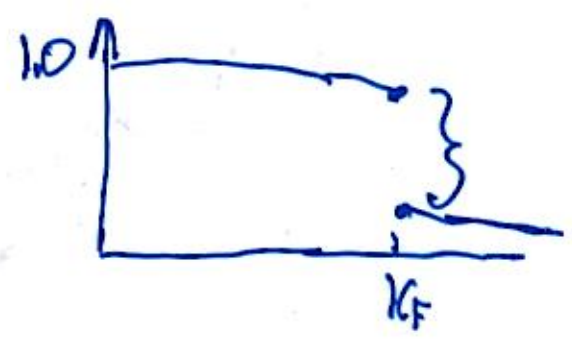
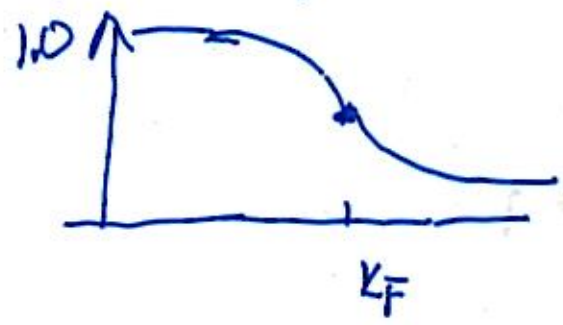
SURPRISE: SCHWINGER ANOMALY

$$[P_1(q), P_1(q')] = \frac{q' L}{2\pi} \delta_{q, -q}$$

↑
 BOSONS



$$\langle a_{1k}^\dagger a_{1k} \rangle = n_{1k}$$



LUTTINGER MODEL
(AT ZERO TEMPERATURE)

LANDAU FERMIL LIQUID

$$\psi_1(x) = \sum_k a_{1k} e^{ikx} c_{1k}$$

↑
↑
 BOSON FERMIONS

THE OTHER WAY AROUND?

$$\psi_1^\dagger(x) = \frac{1}{\sqrt{L}} \sum_k a_{1k}^\dagger e^{-ikx}$$

$$\text{AND } \psi_1(x) = \int e^{2i\phi(x) + \phi_1(x)}$$

$$\phi_1(x) = \frac{2\pi}{L} \sum_{k \neq 0} \frac{p_1(k)}{k} e^{-ikx}$$

L. PESCHEL & A. L.

COLEMAN, MANDELSTAM, MATTLIS

CONNECTIONS

STATISTICAL MECHANICS

1. B-VERTEX MODEL (BAXTER)

FIELD THEORY

2. THIRRING MODEL

3. MASSIVE THIRRING MODEL

SOLID STATE

4. X-Y-Z SPIN $\frac{1}{2}$ CHAIN

BOOKS

TSVELIK, MERSESYAN, STONE, MATTIS

1D HUBBARD MODEL (SPIN)

$$H = \sum_{k, \sigma} (-t \cos ka) c_{k\sigma}^\dagger c_{k\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

(SAME SITE REPULSION)

BOSONIZATION GIVES NEW PHYSICS:

SPIN-CHARGE SEPARATION
(RESCHEL, EMERY, B.L.)

$$H \rightarrow H_{\text{SPIN}} + H_{\text{CHARGE}}$$

(NOT FERMION LIQUID SINCE A QUASI-PARTICLE EXCITATION CARRIES BOTH.)

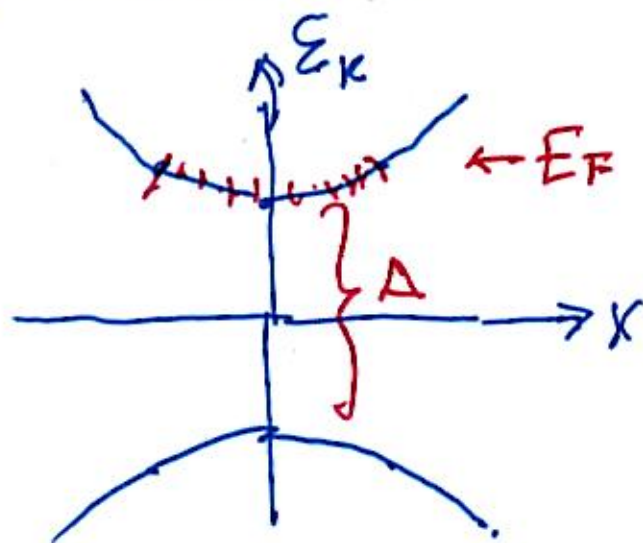
NEED SUSCEPTIBILITIES:

HUTTINGER:

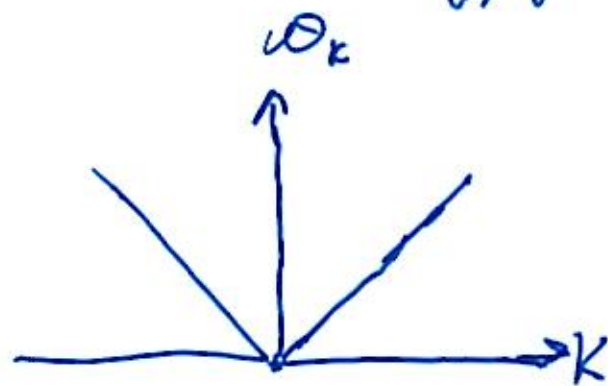
$$\chi_{SG}(T) = \chi_0 \left(\frac{T}{T_0} \right)^p$$

$p > 0$ FOR REPULSION

1D NUMBERED



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CHARGE

SPIN

$E_F = 0$ AT HALF FILLING

$E_F > \frac{\Delta}{2}$ TWO PROCESSES

↳ EXCITATIONS AT E_F (GFT)

$$\chi_{SG} = \gamma_0 \left(\frac{T}{T_0} \right)^{\rho} \quad \rho > 20$$

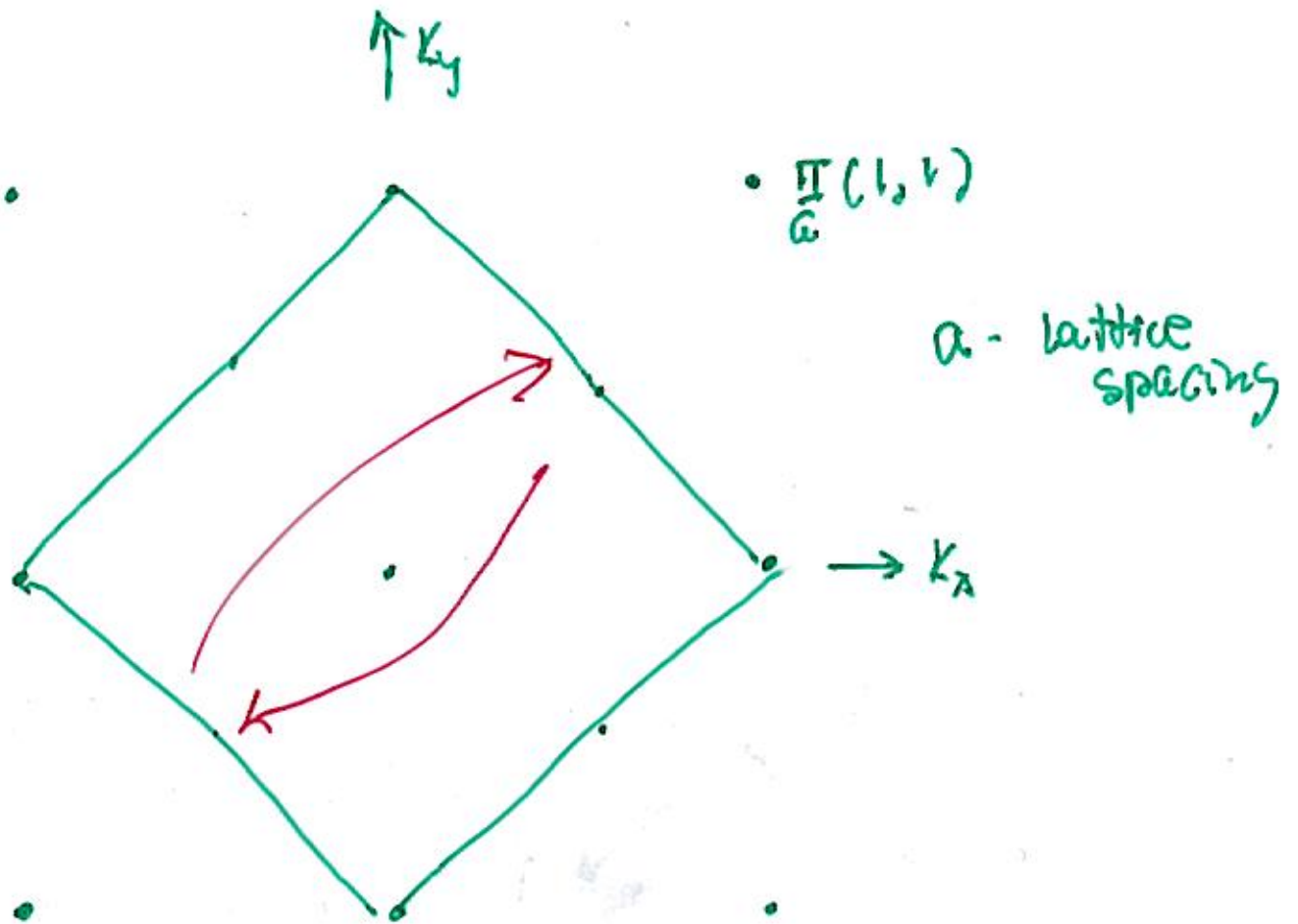
γ_0 PAIR EXCHANGE? $U > 0$

$$\frac{v^2}{(-\Delta)} \quad ?$$

CAN THIS DOMINATE?

2D HUBBARD

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SQUARE FERMI SURFACE, AT HALF FILLING
 (ONE ELECTRON PER SITE)
 NEAREST NEIGHBOR HOPPING

$$H = \sum_{\langle ij \rangle} \sum_{\sigma} t a_{i,\sigma}^{\dagger} a_{j,\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

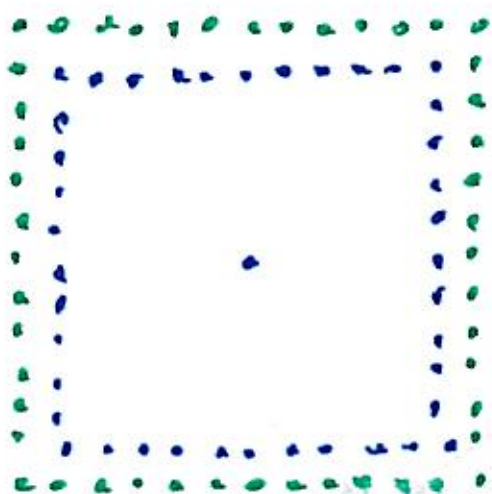
(n.n. sites) (SITES)

SUGGESTED NEW MODEL

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ENERGY SQUARES

↑ k_y



→ k_x

$$E(k_x, k_y) = v k_x \quad (-k_x < k_y < k_x)$$

ALL STATES ON A SQUARE ARE DEGENERATE.

e.g.

$$a_Q(k_x) = \sum_{k_y} e^{iQ k_y} a(k_x, k_y)$$

Q → "ANGULAR" MOMENTUM

AT HALF FILLING → HUBBARD MODEL

LUTTINGER ↔ HUBBARD?

WHY THIS SQUARE ENERGY MODEL?

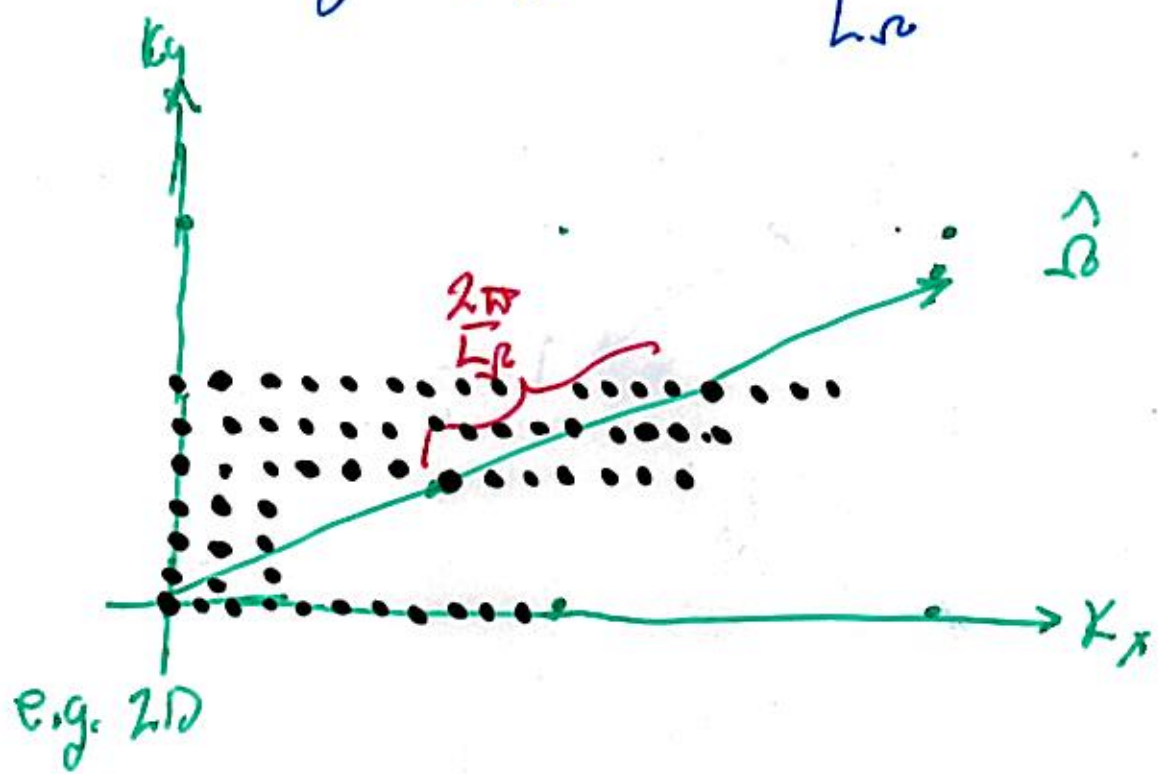
1. CAPTURES PHYSICS OF 2D HUBBARD
2. INTEGRABLE \rightarrow EXACT
3. NEED EXACT TO MAKE BELIEVABLE PREDICTIONS
4. WANT TO FIND CONDITIONS FOR ROOM TEMPERATURE SUPERCONDUCTORS.

IN 3D QED

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$$\vec{A}(\vec{x}) = \frac{1}{L^{3/2}} \sum_{\vec{k}, \lambda} \frac{(b_{\vec{k}, \lambda}^\dagger + b_{-\vec{k}, \lambda})}{\sqrt{2|\vec{k}|}} e^{-i\vec{k} \cdot \vec{x}} \hat{\epsilon}_{\vec{k}, \lambda}$$

$$\vec{k} = \frac{2\pi}{L} (n_x, n_y, n_z) = \frac{2\pi}{L} \vec{n}$$



EQUALLY SPACED POINTS
 (SCHOTTE & A.L.) BOSON ↔ FERMION DUALITY
 CONNECTION TO NUMBER THEORY

$\vec{A}(\vec{x}) \rightarrow$ FERMIONS $b_{\vec{k}}^\dagger \sim \sum_m c_{\vec{n}+\vec{m}, \lambda}^\dagger a_{\vec{m}, \lambda}$
 $\vec{J} \cdot \vec{A}$ HAS ONLY BACKSCATTERING.
 INTEGRABLE.

QED EIGENSTATES IN 3+1

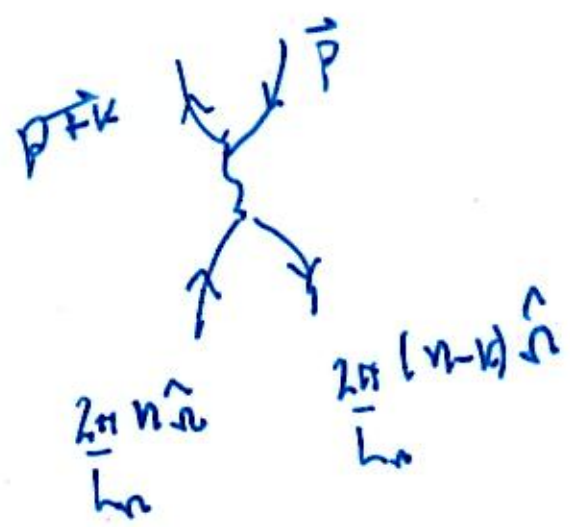
$$H = \sum_{\vec{k}} \bar{\Psi}_{\vec{k}} (\vec{k} \cdot \vec{\alpha} + \beta m) \Psi_{\vec{k}} + e^2 \sum_{\vec{k}} \vec{J}(\vec{k}) \cdot \vec{A}(\vec{k}) + \sum_{\vec{k}, \lambda} \omega_{\vec{k}} b_{\vec{k}, \lambda}^{\dagger} b_{\vec{k}, \lambda}$$

$$\vec{J}(\vec{k}) = \sum_{\vec{p}} \bar{\Psi}_{\vec{p}+\vec{k}} \vec{\sigma} \Psi_{\vec{p}}$$

$$\vec{A}(\vec{k}) = \frac{1}{\sqrt{2k}} \sum_{\lambda} \hat{e}_{\lambda} (b_{\vec{k}, \lambda}^{\dagger} + b_{-\vec{k}, \lambda})$$

$$\vec{k} = \frac{2\pi n \hat{n}}{L} \rightarrow \sum_n \frac{1}{n} C_n^{\dagger} C_{n+k} C_n$$

HAS ONLY BACKSCATTERING



SUMMARY

BOSONS \leftrightarrow FERMIONS

1D: NEED USER FRIENDLY EIGENFUNCTIONS
IN ORDER TO CALCULATE SUSCEPTIBILITIES

2D: THE "SQUARE" MODEL, TO UNDERSTAND
2D HUBBARD MODEL. ARE THERE BIG
SUPERCONDUCTING FLUCTUATIONS?

3D: QED AS A PURELY FERMION BACKSCATTER
PROBLEM. QCD?

SUPERSYMMETRY \leftrightarrow DUALITY

\neq



WHAT ARE THE GENERATORS?
THE GROUP?

WHERE IS MS. KING?

CONNECTION TO NUMBER THEORY

MÖBIUS FUNCTION

$$\sum_{m=1}^{\infty} \mu(m) f(mn) = \delta_{n,1} f(1)$$

m, n integers

$$\mu(m) = \begin{cases} +1 & \text{if } m=1 \\ -1 & \text{if } m \text{ is square-free} \\ 0 & \text{if } m \text{ is not square-free} \end{cases}$$

$$\sum_n e^{\frac{2\pi i \hat{\Omega} \cdot \vec{x}}{L_n}} = ?$$

$$\sum_{\vec{k}} = \sum_n \sum_m \left(\vec{k} = \frac{2\pi n \hat{\Omega}}{L_n} \right)$$

~~$$\sum_n e^{\frac{2\pi i \hat{\Omega} \cdot \vec{x}}{L_n}} = \sum_{m,n} \dots$$~~

$$\begin{aligned} \sum_{m,n} e^{i m \vec{k} \cdot \vec{x}} \mu(m) &= \sum_{m,n} e^{\frac{2\pi i m n \hat{\Omega} \cdot \vec{x}}{L_n}} \mu(m) \\ &= \sum_n e^{\frac{2\pi i \hat{\Omega} \cdot \vec{x}}{L_n}} \\ &= \sum_{\vec{k}, m} \mu(m) e^{i m \vec{k} \cdot \vec{x}} \end{aligned}$$