Quantum Field Theory for Condensed Matter - 2017 Exercise Set 1 (14 points) Due date: monday, april 10th

1. (0.5 p) Show that for a set of operators A, B and C

a)

$$[AB, C] = A[B, C] + [A, C]B.$$
(1)

This can be useful for bosonic systems.

b)

$$[AB, C] = A\{B, C\} - \{A, C\}B.$$
(2)

This can be useful for fermionic systems.

2. (0.5 p) Fermion creation $(c^{\dagger}_{\sigma}(\mathbf{r}))$ and annihilation $(c_{\sigma}(\mathbf{r}))$ operators in real space obey the anti-commutation algebra,

$$\{c_{\sigma}(\mathbf{r}), c_{\sigma'}^{\dagger}(\mathbf{r}')\} = \delta_{\sigma\sigma'}\delta(\mathbf{r} - \mathbf{r}'), \qquad \{c_{\sigma}(\mathbf{r}), c_{\sigma'}(\mathbf{r}')\} = 0.$$
(3)

Assume that the operator $a_{i\sigma}^{\dagger}$ creates an electron in the state $|W_i\rangle$ which is a Wannier state localized at site *i* of the lattice. Using a change of basis and the orthonormality of the Wannier states, show that these operators obey the fermionic algebra as well,

$$\{a_{i\sigma}, a_{j\sigma'}^{\dagger}\} = \delta_{\sigma\sigma'}\delta_{ij}, \qquad \{a_{i\sigma}, a_{j\sigma'}\} = 0.$$

$$\tag{4}$$

3. a) (0.5 p) Consider a general quadratic Hamiltonian

$$H = \sum_{ij} a_i^{\dagger} \mathcal{H}_{ij} a_j, \tag{5}$$

where a_i^{\dagger} could be a boson or fermion creation operator. What is the constraint on \mathcal{H} such that the hamiltonian H is hermitian?

b) (0.5 p) Assume that a_i is either a boson or fermion annihilation operator. If we perform a linear transformation like,

$$d_i = \sum_j M_{ij} a_j,\tag{6}$$

what requirement do we have to put on M in order to preserve the algebra, i.e. the operators d_i obey the same bosonic or fermionic algebra.

Furthermore, show that the discrete Fourier transform,

$$a_{k\sigma} = \frac{1}{\sqrt{\mathcal{N}}} \sum_{i} e^{-i\mathbf{k}.\mathbf{r}_{i}} a_{i\sigma},\tag{7}$$

where \mathcal{N} is the number sites and \mathbf{r}_i is the position of the *i*th lattice point, meets the requirement.

- 4. a) (1.5 p) Take a square lattice of size $N \times N$ with lattice spacing a and set $\mathcal{H}_{ij} = -t$ for nearest neighbours and zero otherwise in the Exercise 3.a) for a fermionic system. Using a discrete Fourier transformation calculate eigenvalues of the Hamiltonian $\epsilon_{\mathbf{k}}$.
 - b) (0.5 p) Sketch the contours of constant $\epsilon_{\mathbf{k}}$ in the Brillouin zone and note the geometry at half-filling (i.e., the Fermi surface obtained by filling exactly half of all available states).

5. (2.0 p) Consider the translational invariant two-body interaction,

$$\hat{V} = \frac{1}{2} \int d^d x d^d y \psi^{\dagger}(\mathbf{x}) \psi^{\dagger}(\mathbf{y}) V(\mathbf{x} - \mathbf{y}) \psi(\mathbf{y}) \psi(\mathbf{x}).$$
(8)

Perform a Fourier transformation and rewrite the two-body interaction \hat{V} in momentum space in terms of

$$\tilde{V}(\mathbf{q}) = \int d^d r e^{-i\mathbf{q}.\mathbf{r}} V(\mathbf{r}).$$
(9)

- 6. Consider a spin chain in which at each site there is a spin-S 'particle'.
 - a) (0.5 p) Recall that the SU(2) algebra is:

$$S_m^+, S_n^-] = 2\delta_{mn} S_m^z, \qquad [S_m^z, S_n^\pm] = \pm \delta_{mn} S_m^\pm, \tag{10}$$

where $S_m^{\pm} = S_m^x \pm i S_m^y$. Show that this algebra can be represented by the following expressions known as the **Holstein-Primakoff** transformation:

$$S_m^+ = (2S - a_m^{\dagger} a_m)^{\frac{1}{2}} a_m, \qquad S_m^- = a_m^{\dagger} (2S - a_m^{\dagger} a_m)^{\frac{1}{2}}, \qquad S_m^z = S - a_m^{\dagger} a_m, \tag{11}$$

in which a_m is a bosonic annihilation operator at site m, obeying the bosonic algebra.

- b) (0.5 p) Using part a, elaborate on the connection between the "harmonic oscillator" states constructed by a^{\dagger} and the spin state at each site. Can this harmonic oscillator accept any occupation number? What are the translations of the states $S_z = -S$ and $S_z = +S$ to the harmonic states?
- c) (1.0 p) Apply the Holstein-Primakoff transformation and a Taylor expansion for large S to the Ferromagnetic Heisenberg hamiltonian and derive the dispersion relation for spin waves in this case.
- 7. Consider a general linear transformation, a so called Bogoliubov transformation,

$$d_i = \sum_j A_{ij} c_j + B_{ij} c_j^{\dagger}.$$

a) (1.0 p) Show that the requirements for d_i to have the same bosonic or fermionic algebra as c_i , *i.e.*,

$$\left[d_i, d_j^{\dagger}\right]_{\pm} = \delta_{ij}, \qquad \left[d_i, d_j\right]_{\pm} = 0,$$

where $[,]_{-}$ denotes a commutator and $[,]_{+}$ denotes an anti-commutator, are the following ones:

$$AA^{\dagger} \pm BB^{\dagger} = 1 \qquad AB^{T} \pm BA^{T} = 0,$$

where + is fermions and - for bosons.

b) (1.0 p) For the fermionic case, show that AB^T is antisymmetric and that the first requirement in part a) can be satisfied by setting

$$A_{ij} = \cos \theta_i U_{ij}, \qquad B_{ij} = \sin \theta_i V_{ij}$$

where U and V are unitary matrices and the θ_i are real.

c) (1.0 p) For the bosonic case, show that AB^T is symmetric and that the first requirement in part a) can be satisfied by setting

$$A_{ij} = \cosh \theta_i U_{ij}, \qquad B_{ij} = \sinh \theta_i V_{ij}$$

where U and V are unitary matrices and the θ_i are real.

8. Take the general quadratic Hamiltonian for spinless fermions involving two sites,

$$H = \sum_{ij=1,2} c_i^{\dagger} h_{ij} c_j + \Delta (c_1 c_2 + c_2^{\dagger} c_1^{\dagger}), \qquad (12)$$

where h_{ij} is a matrix obeying the constraint that you found in Exercise 3. a) and Δ is a real constant.

a) (1.0 p) Using the notation of Exercise. 7, consider the following Bogoliubov transformation:

$$\begin{pmatrix} d_1 \\ d_2 \end{pmatrix} = \begin{pmatrix} \sin\theta & 0 \\ 0 & -\sin\theta \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} + \begin{pmatrix} 0 & \cos\theta \\ \cos\theta & 0 \end{pmatrix} \begin{pmatrix} c_1^{\dagger} \\ c_2^{\dagger} \end{pmatrix}$$
(13)

and show that it satisfies the requirements of Exercise 7. a) for fermions for real θ . In part c) we will determine θ .

- b) (1.0 p) Invert the transformation and write down c_1 and c_2 in terms of $d_1, d_1^{\dagger}, d_2, d_2^{\dagger}$ and θ .
- c) (1.0 p) Insert what you found in part b) into the hamiltonian, and find θ such that coupling in front of terms like d_1d_2 and $d_1^{\dagger}d_2^{\dagger}$ vanish. This means that the hamiltonian can be written as,

$$H = \sum_{ij=1,2} d_i^{\dagger} \tilde{h}_{ij} d_j.$$
⁽¹⁴⁾